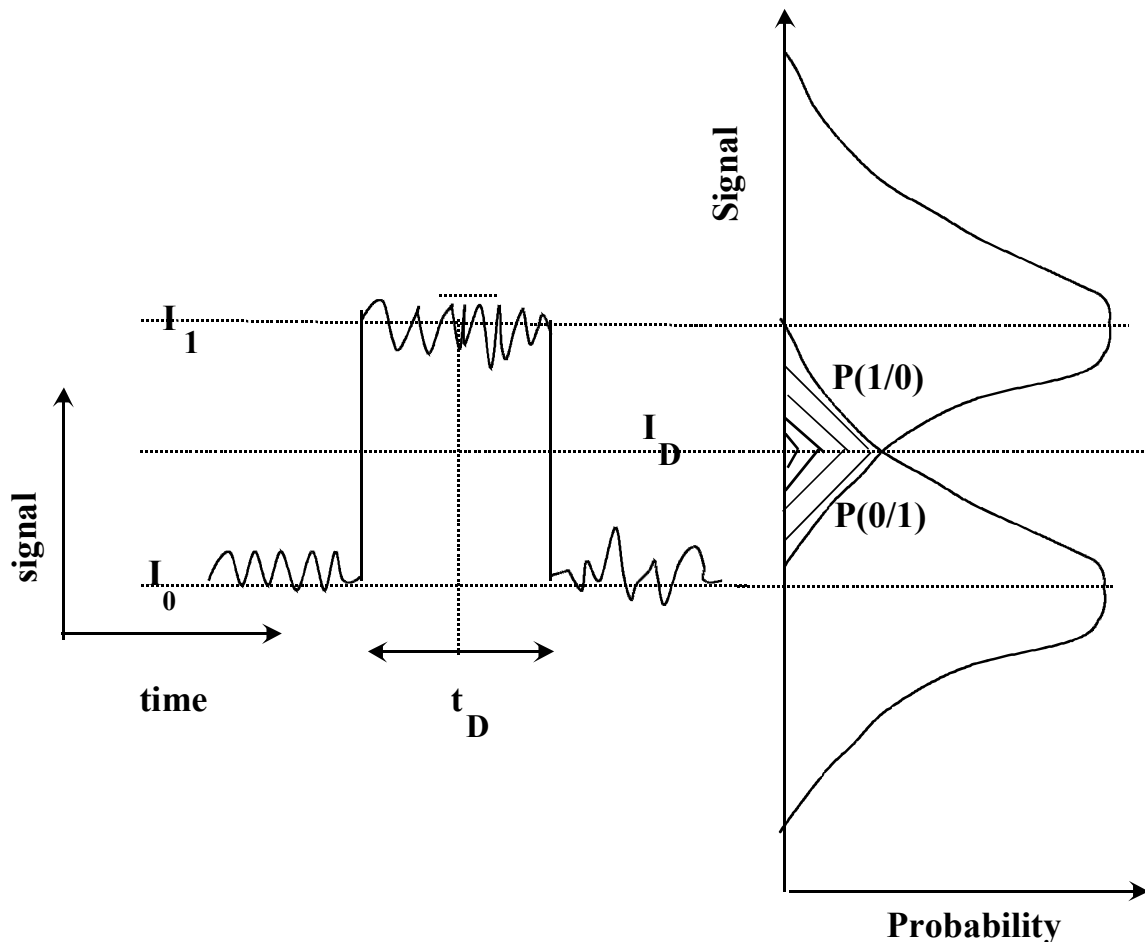


RECEIVER SENSITIVITY:

A measure of a **good receiver** is to have the *same performance* with the *lowest level of incident optical power*.

- **Performance** can be measured as a *low bit error rate (BER)*.
- **BER** \equiv *probability of an incorrect identification of a bit by the decision circuit of a receiver*.
- A BER of 3×10^{-7} implies that 3 bits will be in error for every 10 million received bits.
- For communications systems a criteria used is that the $BER \leq 10^{-9}$.

Receiver Sensitivity \equiv *minimum average received optical power* \bar{P}_{\min} required to achieve a fixed BER.



- If $I > I_D$ a 1 bit is recorded.
- If $I < I_D$ a 0 bit is registered.
- An error occurs for a 1 bit if $I < I_D$ and similarly an error for a zero bit takes place when $I > I_D$.

These errors can be included in the error probability as

$$BER = p(1)P(0/1) + p(0)P(1/0)$$

- $p(1)$ and $p(0)$ are the probabilities of receiving a 1 and 0 bit respectively
- $P(0/1)$ is the probability of deciding 0 when a 1 is received and $P(1/0)$ is the probability of deciding 1 when a 0 is received.
- Assuming that 1s and 0s are equally likely, $p(1) = p(0) = 0.5$

$$BER = 0.5[P(0/1) + P(1/0)].$$

Both *shot* and *thermal noise* can be modeled with *Gaussian statistics* (for large numbers of electrons).

- The total *variance in the current* for I_1 or I_0 can be modeled as

$$\sigma_{1,0}^2 = \langle i_{s,1,0}^2 \rangle + \langle i_{T,1,0}^2 \rangle$$

- The *average noise currents* squared were previously described for shot and thermal noise.
- Since the *average current* I_p is different for the 1 and 0 levels the shot noise and the variance will be different.
- Consequently in general σ_1 and σ_0 will also differ.

- The *conditional probabilities* based on Gaussian statistics are:

$$P(0/1) = \frac{1}{\sigma_1 \sqrt{2\pi}} \int_{-\infty}^{I_D} \exp\left(-\frac{(I - I_1)^2}{2\sigma_1^2}\right) dI = \frac{1}{2} \operatorname{erfc}\left(\frac{I_1 - I_D}{\sigma_1 \sqrt{2}}\right)$$

$$P(1/0) = \frac{1}{\sigma_0 \sqrt{2\pi}} \int_{I_D}^{\infty} \exp\left(-\frac{(I - I_0)^2}{2\sigma_0^2}\right) dI = \frac{1}{2} \operatorname{erfc}\left(\frac{I_D - I_0}{\sigma_0 \sqrt{2}}\right)$$

$$\operatorname{erfc}(y) = \frac{2}{\sqrt{\pi}} \int_x^{\infty} \exp(-y)^2 dy$$

This results in a BER of:

$$\operatorname{BER} = \frac{1}{4} \left[\operatorname{erfc}\left(\frac{I_1 - I_D}{\sigma_1 \sqrt{2}}\right) + \operatorname{erfc}\left(\frac{I_D - I_0}{\sigma_0 \sqrt{2}}\right) \right].$$

I_D is chosen to minimize BER. This occurs when

$$(I_1 - I_D) / \sigma_1 = (I_D - I_0) / \sigma_0 \equiv Q$$

and

$$I_D = \frac{\sigma_0 I_1 + \sigma_1 I_0}{\sigma_0 + \sigma_1}.$$

- When $\sigma_1 = \sigma_0$, $\rightarrow I_D = (I_1 + I_0)/2$

or the *decision threshold current* I_D lies between I_0 and I_1 .

In this case:

$$BER = \frac{1}{2} \operatorname{erfc}\left(\frac{Q}{\sqrt{2}}\right) \approx \frac{\exp(-Q^2/2)}{Q\sqrt{2\pi}}.$$

The approximation is good for $Q > 3$.

Solving for Q :

$$Q = \frac{I_1 - I_0}{\sigma_1 + \sigma_0}$$

- A BER $\approx 10^{-9}$ occurs with $Q \sim 6$.
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Minimum Received Optical Power:

Since the currents I_1 and I_0 are related to the optical power through the responsivity, it is possible to determine the *minimum optical power* to obtain a *fixed BER*.

For the case with $I_0 = 0$, $P_0 = 0$. The detector current for the 1 bit is

$$I_1 = MRP_1 = 2MR\bar{P}_{\min},$$

where \bar{P}_{\min} is the *average received minimum power* with

$$\bar{P}_{\min} = (P_1 + P_0)/2.$$

Noise terms include both *thermal* and *shot noise* for the 1 bit *but only thermal noise* for the 0 bit:

$$\sigma_1 = \left(\sigma_s^2 + \sigma_T^2 \right)^{1/2}; \sigma_0 = \sigma_T .$$

Therefore:

$$Q = \frac{I_1}{\sigma_1 + \sigma_0} = \frac{2MR\bar{P}_{\min}}{\left(\langle i_s^2 \rangle + \langle i_T^2 \rangle \right)^{1/2} + \langle i_T^2 \rangle^{1/2}} .$$

Solving for P_{\min}

$$P_{\min} = \frac{Q}{R} \left(qF_A Q B_e + \frac{\langle i_T^2 \rangle^{1/2}}{M} \right)$$

Quantum Limit for Photodetection:

- An ideal photodiode (no thermal noise, I_D , and 100% QE), has $\sigma_0 = 0$ (no shot noise).
- In this case $I_0 \sim 0$, very few photons (**even 1!**) can be detected provided that they produce *e-h pairs*.
- \therefore In this case *Gaussian statistics* are no longer valid.
- *Poisson statistics* should be applied.

Assuming that N_p is the *average number of photons* in each 1 bit (0 are in a 0 bit) the probability for the formation of m , *e-h hole pairs* is given by *Poisson distribution*

$$P_m = \exp(-N_p) N_p^m / m!$$

- Using our earlier definitions, $P(1/0) = 0$ since no *e-h pairs* are formed when $N_p = 0$.

- $P(0/1) = \exp(-N_p)$

In this case: $m = 0$ since a 0 is decided in this case even though a 1 is received.

$$\therefore BER = \exp(-N_p) / 2.$$

The minimum received optical power in this case is

$$P_{\min} = N_p h \nu B / 2.$$

- At $1.55 \mu\text{m}$ $h\nu = 0.8 \text{ eV}$, and the $P_{\min} = 13 \text{ nW}$ (-48.9 dBm @ 10 Gb/s).
- In most receivers however N_p is typically ~ 1000 which is considerably higher than this limit.

Extinction Ratio:

In many cases the power in the 0-bit will not be 0.

- *Extinction ratio* – can be defined as

$$r = P_o / P_1.$$

- In this case

$$Q = \left(\frac{1-r}{1+r} \right) \frac{2RP_{\min}}{\sigma_1 + \sigma_0}$$

- *Minimum required optical power* becomes

$$P_{\min} = \left(\frac{1+r}{1-r} \right) \cdot \frac{(\sigma_1 + \sigma_0)}{2} \cdot \left(\frac{Q}{R} \right)$$

Effect of Intensity Noise:

- Most optical sources will exhibit some form of *output power fluctuation*.
- This fluctuation is another source of system noise and must be considered in the SNR.
- Assuming that all noise sources are independent the *total current noise power* term can be represented as:

$$\langle i_N^2 \rangle = \langle i_s^2 \rangle + \langle i_T^2 \rangle + \langle i_I^2 \rangle.$$

- The *laser power fluctuation* can be expressed as:

$$\langle i_I \rangle = R \langle (\Delta P_{inc}^2) \rangle^{1/2} = RP_{inc} r_I,$$

with

$$r_I^2 = \frac{1}{2\pi} \int_{-\infty}^{\infty} RIN(\omega) d\omega.$$

RIN(ω) was described previously.

- $r_I = 1/SNR$ for light emitted from the laser source.
- A source with SNR > 20 dB has $r_I < 0.01$.

- With *source noise* the *Q parameter* becomes (with $I_0 = 0$ and $\sigma_0 = \sigma_T$):

$$Q = \frac{I_1 - I_0}{\sigma_1 + \sigma_0} = \frac{2R\bar{P}_{rec}}{\left(\langle\sigma_T^2\rangle + \langle\sigma_s^2\rangle + \langle\sigma_I^2\rangle\right)^{1/2} + \sigma_T}$$

with

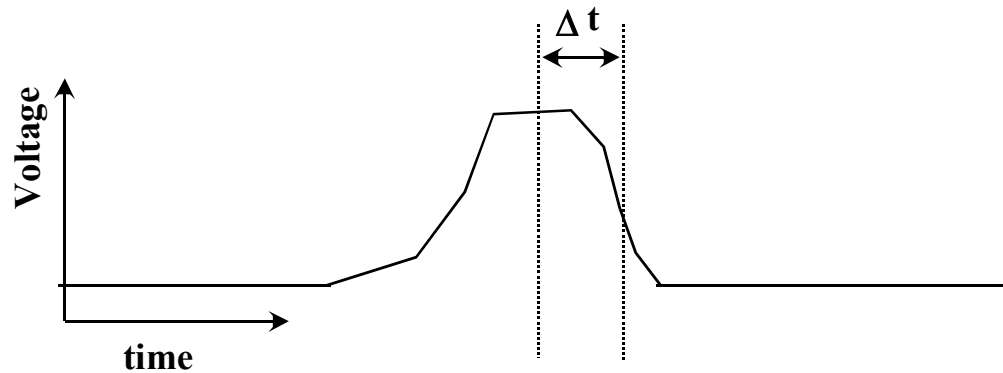
$$\sigma_s = (4qR\bar{P}_{rec}B_e)^5; \sigma_I = 2r_I R\bar{P}_{rec}$$

- The *receiver sensitivity* becomes:

$$\bar{P}_{rec}(r_I) = \frac{Q\sigma_T + Q^2qB_e}{R(1 - r_I^2Q^2)}$$

Effects of Timing Jitter:

- The previous receiver sensitivity analysis is based on the assumption that the *signal is sampled at the peak of the voltage pulse*.
- Timing decisions are typically determined by a *clock-recovery circuit*.
- Since the input to this circuit is noisy the *sampling time fluctuates* from bit to bit.
- This fluctuation is referred to as *timing jitter* and will degrade the SNR.



- Since Δt is a *random variable* the *change in the sampled voltage level* is also a *random variable*.
- This reduces the receiver performance.
- The SNR can be improved at the expense of higher source power.

Effect on System Performance:

- Assume a *pin* detector/receiver that is limited by thermal noise with (extinction ratio) $r = 0$, and $I_0 = 0$.

$$Q = \frac{I_1 - \langle \Delta i_j \rangle}{(\sigma_T^2 + \sigma_j^2)^5 + \sigma_T}$$

- assuming that the $B\Delta t \ll 1$ where $B = \text{bit rate} = 1/T_B$,

$$\Delta i_j = (2\pi^2 / 3 - 4) \cdot (B\Delta t)^2 I_1.$$

Defining

$$b = (4\pi^2 / 3 - 8) (B\tau_j)^2$$

with τ_j the RMS value of Δt , the receiver sensitivity becomes

$$\bar{P}_{rec}(b) = \left(\frac{\sigma_T Q}{R} \right) \frac{1 - b/2}{(1 - b/2)^2 - b^2 Q^2 / 2}.$$

- A *power penalty* can be defined as the *increase in received power required to maintain the same BER as obtained without a timing jitter factor*.

$$\delta_j = 10 \log_{10} \left(\frac{\bar{P}_{rec}(b)}{\bar{P}_{rec}(0)} \right) = 10 \log_{10} \left(\frac{1 - b/2}{(1 - b/2)^2 - b^2 Q^2 / 2} \right).$$

The *power penalty* is negligible when *timing jitter* $\tau_j < (0.1)T_B$ but increases rapidly beyond this point.